

Optimal Control of SonoVue Microbubbles to Estimate Hydrostatic Pressure

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Abstract—The measurement of cardiac and aortic pressures enable diagnostic insight into cardiac contractility and stiffness. However, these pressures are currently assessed invasively using pressure catheters. It may be possible to estimate these pressures less invasively by applying microbubble ultrasound contrast agents as pressure sensors. The aim of this study was to investigate the subharmonic response of the microbubble ultrasound contrast agent SonoVue (Bracco Spa, Milan, Italy) at physiological pressures using a static pressure phantom. A commercially available cell culture cassette with Luer connections was used as a static pressure chamber. SonoVue was added to the phantom, and radiofrequency data were recorded on the Ultrasound Advanced Open Platform (ULA-OP). The mean subharmonic amplitude over a 40% bandwidth was extracted at 0–200 mmHg hydrostatic pressures, across 1.7–7.0 MHz transmit frequencies and 3.5–100% maximum scanner acoustic output. The Rayleigh-Plesset equation for single bubble oscillations and additional hysteresis experiments were used to provide insight into the mechanisms underlying the subharmonic-pressure response of SonoVue. The subharmonic amplitude of SonoVue increased with hydrostatic pressure up to 50 mmHg across all transmit frequencies, and decreased thereafter. A decreasing microbubble surface tension may drive the initial increase in the subharmonic amplitude of SonoVue with hydrostatic pressure, while shell buckling and microbubble destruction may contribute to the subsequent decrease above 125 mmHg pressure. In conclusion, a practical operating regime that may be applied to estimate cardiac and aortic blood pressures from the subharmonic signal of SonoVue has been identified.

Index Terms—ultrasound contrast agents, subharmonic imaging, hydrostatic pressure.

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I. INTRODUCTION

MICROBUBBLE based ultrasound contrast agents are currently used in the clinic to complement standard B-mode imaging across multiple organs and systems in the human body, including the heart, breast and liver [1]–[4]. In the heart, the commercially available contrast agents SonoVue/Lumason (Bracco Spa, Milan, Italy), Luminity/Definity (Lantheus Medical Imaging Inc., N. Billerica, MA, USA) and Optison (GE Healthcare, Princeton, NJ, USA) may be applied to assess left ventricular function, structural left ventricular abnormalities and myocardial perfusion [2], [3]. In addition, these microbubbles have a promising new application as pressure sensors [5], [6], which would enable minimally invasive estimations of the cardiac and large artery pressures underpinning diagnostic information on cardiac contractility and stiffness [7]–[9]. Successful implementation of microbubble based cardiac pressures would provide a safer and more cost-effective alternative to the current clinical method requiring invasive cardiac catheterization [8], [9].

Using single element transducers, the ultrasound contrast agents Levovist, Optison, Definity, ZFX and Sonazoid have been found to generate a subharmonic signal that is linearly and negatively correlated to static hydrostatic pressures from 0–186 mmHg [10]. This phenomenon may be due to changes in bubble surface tension and shell buckling [11], [12], and is additionally affected by ultrasound settings such as transmit frequency, pulse length and acoustic pressure. As acoustic pressure increases, the subharmonic signal of ultrasound contrast agents can be delineated into three distinct phases: (i) occurrence, (ii) growth and (iii) saturation (for an example, see Fig. 3) [13], [14]. In the occurrence and saturation phases, the subharmonic amplitude is stable despite increases in acoustic pressure. In contrast, the growth phase is characterized by an increase in subharmonic amplitude with acoustic pressure. An acoustic pressure within the growth phase has been shown to be necessary to elicit a strong negative linear relationship between the subharmonic amplitude of ultrasound contrast agents and hydrostatic pressure [6], [13].

Among the commercially-available ultrasound contrast agents, Sonazoid has been investigated the most extensively due to its greatest sensitivity to hydrostatic pressure following the study by Halldorsdottir and colleagues [10]. However, Sonazoid is currently not approved nor marketed in Europe [15]. A potential alternative to Sonazoid is the ultrasound contrast agent SonoVue marketed by Bracco Spa, which is

widely used in Europe [3]. To our knowledge, there have been three investigations of SonoVue as a potential pressure sensor [16]–[18], and none of these have been done by the research group spearheading current efforts with Sonazoid. The first study was by Andersen and Jensen [16], who in fact examined the ratio of the subharmonic to fundamental components, instead of the subharmonic as described by Halldorsdottir and colleagues [10]. The second study by Sun and colleagues [18] found an increase in the subharmonic signal with hydrostatic pressure when excited by a 1.33 MHz ultrasound pulse at an acoustic pressure of 300 kPa, but a decrease when excited by a 4 MHz pulse at a similar acoustic pressure (300 kPa). In contrast, the third study by Li and colleagues (which was published in 2018 while data collection for this current study was ongoing) [17] reported a *decrease* in the subharmonic signal with hydrostatic pressure when excited by a 1.33 MHz ultrasound pulse at an acoustic pressure of 350 kPa, but an *increase* from 0–50 mmHg hydrostatic pressure at 4 MHz and the same acoustic pressure (350 kPa), followed by a decrease from 50–180 mmHg. At 4 MHz and 450–500 kPa acoustic pressures, Li and colleagues [17] found a decrease in the subharmonic signal with hydrostatic pressure, as would be expected from previous work with other ultrasound contrast agents [10]. Of specific relevance to the investigation of SonoVue as a potential pressure sensor is the work by Frinking and colleagues that examined SonoVue-like microbubbles from Bracco Research [11] — they found an increase in the subharmonic signal of SonoVue-like microbubbles to hydrostatic pressure at 4 MHz and 50 kPa acoustic pressure, no change in the subharmonic signal at the same transmit frequency (4 MHz) and 200 kPa acoustic pressure, but a decrease at 400 kPa acoustic pressure. The conflicting data in the literature mean that the subharmonic response of SonoVue to hydrostatic pressure has not been established and is still unclear. Building upon the extensive work led by Forsberg and colleagues towards developing Sonazoid as a pressure sensor, it is probable that an experimental protocol similar to that being used with Sonazoid may help clarify the subharmonic response of SonoVue to hydrostatic pressure.

The aim of this study was to investigate the subharmonic response of the ultrasound contrast agent SonoVue at physiological pressures using a static pressure phantom. We hypothesized that the subharmonic signal of SonoVue would exhibit (i) a growth phase with increasing acoustic pressures, and (ii) a negative linear relationship with hydrostatic pressure at an acoustic pressure within this growth phase. The Rayleigh-Plesset equation for single bubble oscillations, combined with an effective bubble surface tension [19], and additional hysteresis experiments were used to provide insight into the mechanisms underlying the empirically-observed subharmonic-pressure response. Part of this work, limited to the data at transmit frequency 5 MHz, was first presented at the 2017 IEEE International Ultrasonics Symposium [20].

II. MATERIALS AND METHODS

A. Static pressure chamber

A phantom capable of maintaining 0–200 mmHg static hydrostatic pressures was developed using a cell culture cassette

with Luer connections (CLINicell 25, 175 μm membrane, 10 mL volume, 6.8 cm \times 3.9 cm \times 3.7 mm, Mabio International, Tourcoing, France), and submerged in a water bath (Fig. 1). Luer connections ensured that the CLINicell chamber was air-tight to maintain stable hydrostatic pressures. A similar cell culture cassette albeit with a thinner membrane (50 μm membrane) has been recently demonstrated as a viable chamber for microbubble studies [21]. The ultrasound transducer was positioned at a 45° angle relative to the cell culture cassette [22]. This enabled a clear region of interest with minimal backscatter from the cassette windows, and concomitantly increased the effective depth of the pressure chamber on the ultrasound image to 5.2 mm. A 1.5 mm magnetic stirrer was inserted into the cassette to maintain a homogenous concentration of microbubbles within the pressure chamber.

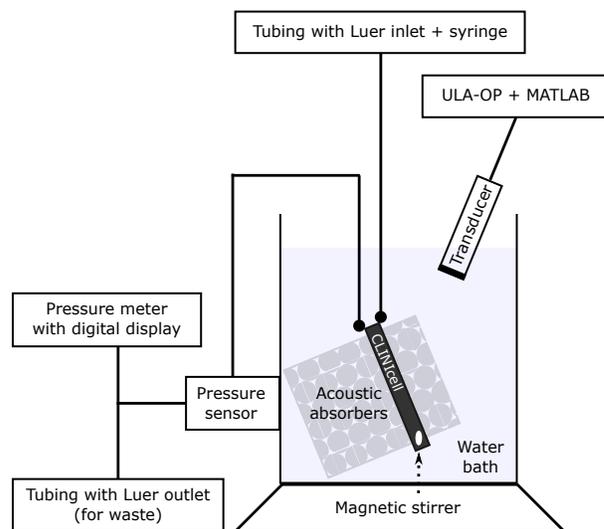


Fig. 1. Schematic representation of the static pressure phantom. The CLINicell was submerged in a water bath with layers of acoustically-absorbent foam in front of and behind it. The transducer was positioned at a 45° angle relative to the window of the CLINicell. SonoVue was added to the CLINicell with a syringe, via a Luer stopcock.

High pressure PVC tubing was secured to the cassette via the two Luer ports (900 PSI, Cole-Parmer, Cambridgeshire, UK), and led out of the water bath to entry and exit Luer stopcocks for administering microbubble solution. Prior to the exit stopcock, a pressure sensor (PRESS-S-000 sensor, PendoTech, Princeton, NJ, USA) was connected and positioned outside the water bath at the same height as the middle of the submerged cassette. The pressure sensor was connected to a digital pressure meter (INFCS-112B meter, Newport Electronics, Inc., Santa Ana, CA, USA) calibrated at 0 mmHg (ambient pressure) and 147 mmHg using a water column (2 m water column).

B. Attenuation and insertion loss

A 12-mm layer of open-cell melamine foam (Basotect, BASF, Ludwigshafen, Germany) was positioned in front of the cassette to create an attenuating layer between the ultrasound transducer and the microbubbles. To measure the attenuation

resulting from this foam and the insertion loss through one window of the cell culture cassette, a pair of broadband transducers (Panametrics V311, 12.7 mm diameter, 59 mm focal length, 10 MHz center frequency; Panametrics V310, 6.35 mm diameter, 5 MHz center frequency; Olympus NDT, Waltham, MA, USA) were used to acquire the through-transmission spectrum using a broadband substitution technique [23]. Measurements were conducted in an 8 L acrylic tank ($45 \times 12 \times 15$ cm) filled with distilled water. An ultrasound pulser-receiver (DPR300, JSR Ultrasonics, Pittsford, NY, USA) was used to generate the excitation pulse and amplify the received signal (20–50 dB gain). Received waveforms were averaged (typically 64 traces/acquisition), digitized (LT264, LeCroy, Chestnut Ridge, NY, USA), and transferred to a computer for analysis using MATLAB (The MathWorks, Inc., Natick, MA, USA). Attenuation through the open-cell melamine foam and insertion loss through one window of the CLINicell were used to calculate total acoustic signal loss, and to estimate the incident acoustic pressure within the cassette chamber (acoustic pressure in a water bath $\times 10^{-\text{total signal loss}/20}$). All tables and figures show peak-negative acoustic pressure corrected for signal loss, unless stated otherwise.

Layers of acoustically-absorbent open-cell foam were additionally positioned behind the cassette to reduce artifacts created due to reflections and scattering beyond the cassette chamber.

C. Experiments with SonoVue at 0–200 mmHg hydrostatic pressures

SonoVue was reconstituted according to manufacturer's instructions, and diluted in gas-equilibrated water to yield the typical concentration used in the clinic (0.4 $\mu\text{L}/\text{mL}$ water). With the exit port open and the magnetic stirrer spinning in the cassette chamber, approximately 25 mL of diluted microbubble solution was added to the static pressure phantom (≈ 0.5 mL/s). The exit port was then closed and hydrostatic pressure was increased by adding more microbubble solution.

Radiofrequency data were recorded across the bandwidth of a linear and a phased array ultrasound transducer (bandwidth 3–7 MHz, LA332E Marzo 2014 and bandwidth 1.2–2.1 MHz, PA230, respectively; Esaote, Genoa, Italy) on the ULtrasound Advanced Open Platform (ULA-OP, MSD Lab, University of Florence, Florence, Italy). These transducers were used in this study to encompass vascular and cardiac imaging.

Pulse-inversion sequences and long transmit pulses were used to enhance the non-linear microbubble signal [18], [24]. A longer transmit pulse has been found to enhance the subharmonic signal-to-noise ratio, and reduce the transient effect of pulse length due to the growth and decay at the beginning and end of each ultrasound pulse on the subharmonic signal [18], [24]. Therefore, pulse length was chosen to maximize the number of cycles per pulse, while not exceeding the 5.2 mm effective chamber depth (with a 45° angle of insonation). 16-cycle pulses were used at transmit frequencies 5–7 MHz, 12 cycle pulses at 4 MHz, 10 cycle pulses at 3 MHz, 7 cycle pulses at 2.1 MHz, and 5 cycle pulses at 1.7 MHz. Pulse length at transmit frequency 5 MHz with 16-cycle pulses was

4.14 mm (standard deviation 0.02 mm, $n = 3$, 82% maximum scanner acoustic output). The maximum mechanical index on the linear transducer was 0.52 at transmit frequency 6 MHz with 16-cycle pulses, and that on the phased array transducer was 0.41 at transmit frequency 2.1 MHz with 7 cycle pulses, as measured with a 0.5 mm hydrophone (SN1832, Precision Acoustics Ltd, Dorchester, UK) in a water bath.

Data were first recorded at ambient hydrostatic pressure (0 mmHg) from 3.5–100% maximum scanner acoustic output ($n = 40$, equally spaced on a logarithmic scale; 9 min per dataset of incremental acoustic outputs) [13] — to determine the acoustic pressure range that elicited the growth phase response of SonoVue. Subsequently, data were recorded across scanner acoustic output levels corresponding to the growth phase ($n = 20$; 4 min per dataset).

Experiments were performed from 200 to 0 mmHg hydrostatic pressures in 25 mmHg decrements, and then repeated. The microbubble solution in the phantom was replenished after each set of acoustic output levels. Experiments at individual hydrostatic pressure levels were further repeated if no crossover was observed between the first two sets (i.e., if one set of data points was consistently higher than the other set within the growth phase). The erroneous dataset was determined based on its large variation from the other two runs that had intersecting growth phases, attributed to human error in replenishing the microbubble solution or noise [25], and discarded. Across the transmit frequencies investigated, additional runs were recorded at 2–6 pressure levels (out of 9) to obtain intersecting growth phases (i.e., 1.7 MHz: 0, 50, 75, 100, 125 and 150 mmHg; 2.1 MHz: 175 and 200 mmHg; 3 MHz: 0, 25, 100, 125, 150 and 175 mmHg; 4 MHz: 0, 25, 150, 175 and 200 mmHg; 5 MHz: 0, 25, 100 and 175 mmHg; 6 MHz: 0, 25, 50, 150, 175 and 200 mmHg; 7 MHz: 0, 175 and 200 mmHg). Noisy subharmonic data is common, and can be mitigated *in vivo* by applying a median filter on a larger number of frames [26]. All experiments were performed at room temperature ($\approx 21^\circ\text{C}$), and were completed within 7 h of microbubble reconstitution. Crossover of the data between repeats indicated that the subharmonic signal of SonoVue was stable across this period.

D. Extracting and analyzing the subharmonic signal

The mean signal amplitude over a 40% bandwidth around the nominal subharmonic frequency (i.e., transmit frequency $f_0/2$) was extracted offline using MATLAB [26]. A zero-phase digital filter was applied with a finite impulse response (FIR) band-pass filter to isolate the signal over the subharmonic bandwidth. In experiments with the linear transducer, a 15.5×3 mm region of interest was defined based on the B-mode image at the center frequency of 5 MHz with 16-cycle pulses (Fig. 2). In experiments with the phased array transducer, a sector region of interest with 3 mm depth was defined from the B-mode image at transmit frequency 2.1 MHz with 7 cycle pulses. The average subharmonic amplitude was calculated as the mean amplitude across three frames of the region of interest (i.e., 3 frames \times 64 lines/frame).

Linear regressions between subharmonic amplitude and hydrostatic pressure were performed at each acoustic output level

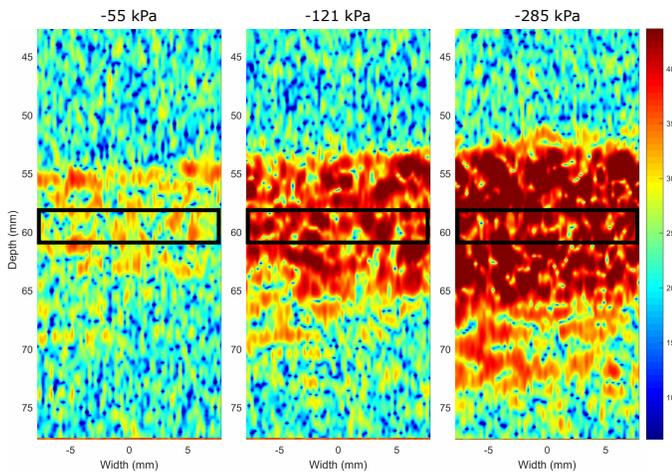


Fig. 2. 2D images of the subharmonic signal at (A) 55 kPa, (B) 121 kPa, and (C) 285 kPa peak-negative acoustic pressures, illustrating a greater amplitude (dB) with increasing peak-negative acoustic pressure. In this example, hydrostatic pressure was maintained at 75 mmHg, and data were recorded at transmit frequency 4 MHz with 12 cycle pulses. The region of interest is demarcated by the black rectangular border.

to identify the maximum sensitivity of SonoVue to changes in hydrostatic pressure. Mean error was calculated as the mean absolute difference between the data and the regression line.

E. Using the Rayleigh-Plesset equation to investigate the subharmonic-pressure relationship of SonoVue

To gain first insights into the mechanisms underlying the subharmonic-pressure relationship of SonoVue, we used a classical model for single bubble oscillations described by the Rayleigh-Plesset equation and assumed an adiabatic gas response. The Peclet number for a two-micron air bubble at 1 MHz is 120, and therefore the adiabatic assumption ($Pe \gg 1$) is verified [27]. This provided the amplitude of bubble oscillation as a function of hydrostatic pressure when the bubble is excited with a pure sinusoidal wave of known frequency f and fixed peak-to-peak acoustic pressure of 150 kPa. The equilibrium bubble radius R_0 at a given pressure p was obtained by assuming that the amount of gas inside the bubble remains constant:

$$R_0^3 \left(p + \frac{2\sigma}{R_0} \right) = C_{\text{ref}} \quad (1)$$

where C_{ref} is a constant obtained from the reference radius R_{ref} measured at the reference pressure. $R_{\text{ref}} = 2 \mu\text{m}$, using the typical modal radius of SonoVue at reference pressure [28].

The amplitude of SonoVue bubble oscillations is additionally affected by buckling of its phospholipid monolayer [19], [28]. As hydrostatic pressure increases, bubble radius decreases and buckling occurs below a critical radius, which is dependent on the number of phospholipid molecules surrounding the bubble. Effective surface tension σ of the lipid monolayer ranges from 0.07 N/m in the elastic state for air/water systems, to 0 N/m in the buckled state (Eq. 1). A parametric study was performed to investigate the influence of surface tension on the amplitude of bubble oscillations.

We then used the predicted oscillation amplitude from the single bubble model to estimate reflected sound as a function of hydrostatic pressure. In the weakly non-linear regime, it is reasonable to assume that the intensity of reflected sound is proportional to the intensity of non-linearity [29] and inversely proportional to the void fraction [30]. These parameters are represented in Equation 2 as bubble oscillation amplitude and equilibrium pressure (p_0), respectively. If we adjust the model using the signal intensity at the reference pressure for air water bubbles, the predicted amplitude A in dB across 0–200 mmHg equilibrium pressures p_0 is:

$$A = 27 + 20 \log \left(\frac{\Delta R_0}{\Delta R_{\text{ref}}} \frac{p_{\text{ref}}}{p_0} \right) \quad (2)$$

F. Hysteresis experiments to investigate the irreversible impact of hydrostatic pressure on SonoVue

SonoVue was added to the static pressure phantom following the experimental protocol in Section II-C. A transmit frequency of 5 MHz with 16-cycle pulses was used for this set of experiments. Similar to the experiments in Section II-C, radiofrequency data were first recorded at ambient pressure (0 mmHg) to determine the range of scanner acoustic output levels corresponding to the growth phase ($n = 40$).

To investigate the irreversible impact of hydrostatic pressure on SonoVue, the microbubble solution in the phantom was first maintained at 200–0 mmHg without insonation. The pressure chamber was then returned to 0 mmHg, and data were recorded at scanner acoustic output levels corresponding to the growth phase ($n = 20$). SonoVue solution in the phantom was replenished after each set of acoustic output levels. Experiments were performed in 25 mmHg decrements. Three sets of data were recorded for each hydrostatic pressure level, and experiments were completed within 4.5 h of microbubble reconstitution.

III. RESULTS

A. Acoustic signal loss before reaching the microbubbles

At the 45° angle of insonation used in this study, total attenuation through the foam and single cassette window ranged from 4.6–16.2 dB across 1.7–7.0 MHz transmit frequencies (Tables I–II). The insertion loss across 0–50° angles of insonation are reported in the Appendix.

TABLE I
ATTENUATION THROUGH OPEN-CELL MELAMINE FOAM (BASOTECT, BASF, GERMANY) ACROSS 1–7 MHz TRANSMIT FREQUENCIES. VALUES ARE MEAN (STANDARD DEVIATION) OF TWO MEASUREMENTS.

Frequency (MHz)	Attenuation (dB/cm)
1.0	0.89 (0.44)
2.0	1.00 (0.49)
3.0	1.12 (0.61)
4.0	1.35 (0.66)
5.0	1.75 (0.60)
6.0	2.15 (0.55)
7.0	2.51 (0.50)

TABLE II
INSERTION LOSS THROUGH ONE WINDOW OF THE CLINICELL CHAMBER AT A 45° ANGLE OF INSONATION, ACROSS 1–7 MHz TRANSMIT FREQUENCIES. VALUES ARE MEAN (STANDARD DEVIATION) OF FOUR MEASUREMENTS.

Frequency (MHz)	Insertion loss (dB)
1.0	0.92 (0.05)
1.5	2.04 (0.06)
2.0	4.36 (0.08)
2.5	10.45 (0.28)
3.0	14.30 (0.39)
3.5	4.34 (0.11)
4.0	2.63 (0.05)
4.5	2.72 (0.09)
5.0	3.22 (0.17)
5.5	3.64 (0.14)
6.0	4.25 (0.03)
6.5	7.23 (0.16)
7.0	10.50 (0.57)

B. Occurrence, growth and saturation phases of SonoVue with increasing acoustic pressure

SonoVue generated subharmonic signals that followed the characteristic occurrence, growth and saturation phases previously observed in microbubble ultrasound contrast agents [13] (representative example in Fig. 3). The acoustic pressures corresponding to these phases varied across transmit frequencies, with the growth phase occurring between 50–250 kPa peak-negative pressure.

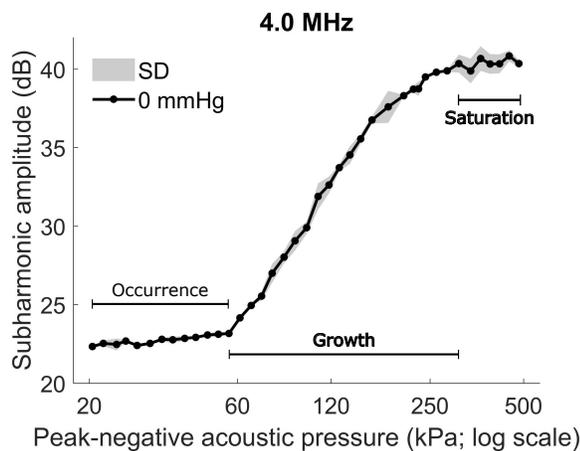


Fig. 3. Mean subharmonic amplitude of SonoVue at ambient hydrostatic pressure (0 mmHg) across the full range of scanner acoustic pressures at transmit frequency 4 MHz ($n = 2$). SD: 1 standard deviation.

C. Ascending and descending phases of the subharmonic amplitude of SonoVue with hydrostatic pressure

From 0 to 50 mmHg hydrostatic pressure, the subharmonic signal of SonoVue increased for all transmit frequencies investigated in this study (1.7–7.0 MHz; Fig. 4). Between 50–75 mmHg hydrostatic pressure, the subharmonic signal further increased for most transmit frequencies except at 1.7 and 3.0 MHz, which corresponded to the lowest transmit frequencies investigated for both the linear and phased array

transducers. For these transmit frequencies, the subharmonic signal plateaued between 50–75 mmHg.

Above 75 mmHg, the subharmonic signal decreased as hydrostatic pressure increased for transmit frequencies 1.7–3.0 MHz. For transmit frequencies 4.0–7.0 MHz, however, the subharmonic signal plateaued between 75–125 mmHg hydrostatic pressure, and only decreased as hydrostatic pressures were increased above 125 mmHg. The maximum sensitivity of the subharmonic-pressure relationship for each transmit configuration was identified for 0–75 mmHg and 125–200 mmHg hydrostatic pressures separately, and summarized in Tables III–IV, and Fig. 5.

TABLE III
MAXIMUM SENSITIVITY OF THE SUBHARMONIC AMPLITUDE OF SONOVUE TO 0–75 MMHG HYDROSTATIC PRESSURE ACROSS 1.7–7.0 MHz TRANSMIT FREQUENCIES

Linear regression results	1.7 MHz	2.1 MHz	3.0 MHz	4.0 MHz	5.0 MHz	6.0 MHz	7.0 MHz
Maximum sensitivity (dB/mmHg)	0.16	0.15	0.13	0.15	0.13	0.07	0.09
Mean error (mmHg)	13.9	4.4	7.6	2.0	1.4	5.6	9.6
Peak-negative pressure (kPa)	126	155	53	121	141	153	140
Mechanical index	0.16	0.24	0.20	0.11	0.13	0.16	0.29
Adjusted r^2	0.677	0.957	0.894	0.992	0.995	0.938	0.820
p	0.007	<0.001	<0.001	<0.001	<0.001	<0.001	0.001

TABLE IV
MAXIMUM SENSITIVITY OF THE SUBHARMONIC AMPLITUDE OF SONOVUE TO 125–200 MMHG HYDROSTATIC PRESSURE ACROSS 1.7–7.0 MHz TRANSMIT FREQUENCIES

Linear regression results	1.7 MHz	2.1 MHz	3.0 MHz	4.0 MHz	5.0 MHz	6.0 MHz	7.0 MHz
Maximum sensitivity (dB/mmHg)	-0.09	-0.09	-0.07	-0.14	-0.16	-0.13	-0.05
Mean error (mmHg)	9.2	7.1	10.9	2.6	3.6	6.9	8.5
Peak-negative pressure (kPa)	233	155	78	101	131	273	191
Mechanical index	0.30	0.24	0.29	0.09	0.12	0.28	0.40
Adjusted r^2	0.871	0.911	0.811	0.986	0.974	0.910	0.864
p	<0.001	<0.001	0.001	<0.001	<0.001	<0.001	<0.001

D. Effects of hydrostatic pressure and surface tension using the Rayleigh-Plesset equation

At a constant surface tension ($\sigma = 0.07$ N/m), the simulations revealed a decrease in subharmonic amplitude with increasing hydrostatic pressure (Fig. 6A). In contrast, at a constant hydrostatic pressure ($p = 0$ mmHg), subharmonic amplitude increased with decreasing surface tension (Fig. 6B). The predicted subharmonic amplitude in dB as a function of hydrostatic pressure and surface tension is shown in Fig. 7.

E. Effects of prior exposure to 0–200 mmHg hydrostatic pressures on the subharmonic signal of SonoVue

The subharmonic amplitude of SonoVue was lower after exposure to 150–200 mmHg hydrostatic pressures, but not

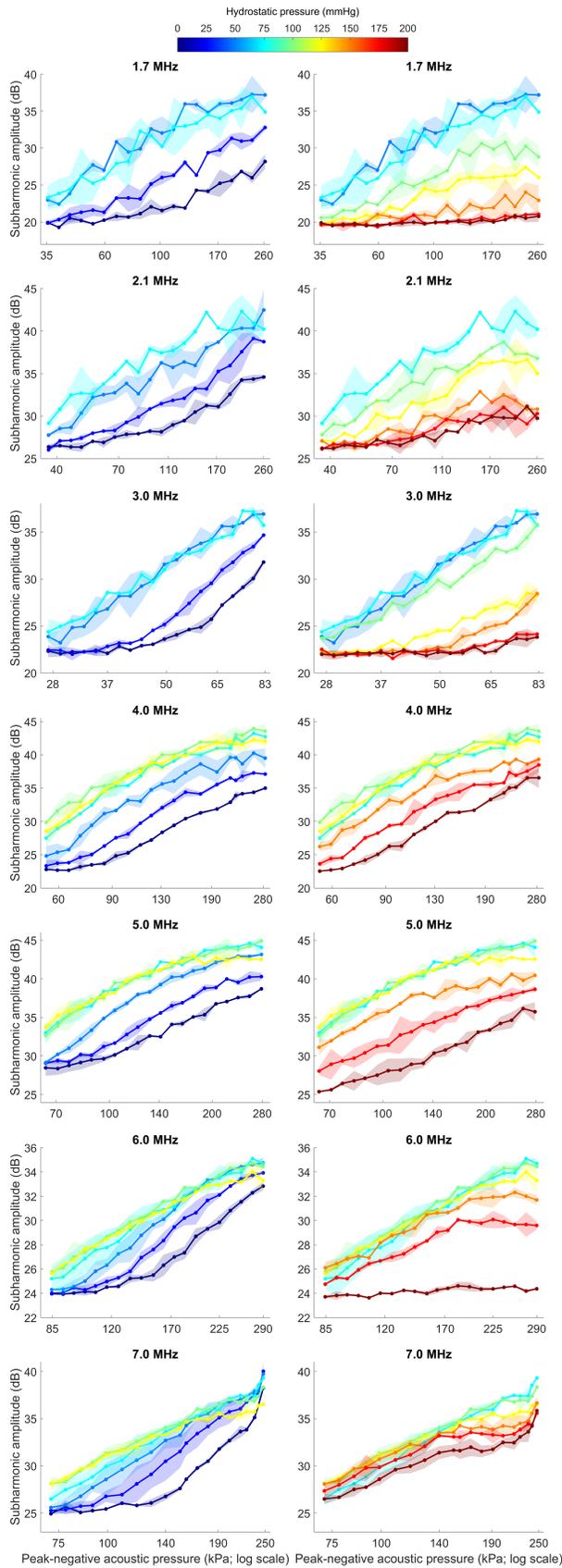


Fig. 4. Mean subharmonic amplitude of SonoVue at 0–200 mmHg hydrostatic pressures, across transmit frequencies 1.7 and 2.1 MHz with the phased array transducer, and 3.0–7.0 MHz with the linear transducer ($n = 2$). The left column shows the ascending phase of the subharmonic-pressure relationship, while the right column shows the descending phase. Data at the plateau phase are repeated in both columns for reference. Translucent shading indicates 1 standard deviation around the mean. The horizontal axis differs between rows for clarity of the individual plots.

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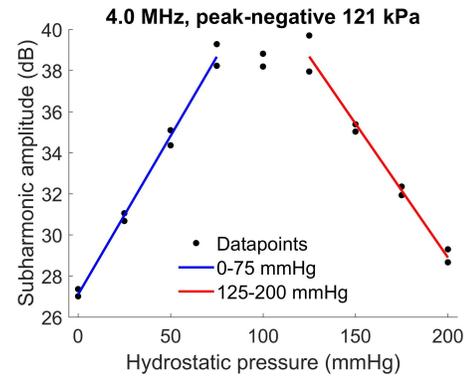


Fig. 5. Subharmonic amplitude of SonoVue from 0–200 mmHg hydrostatic pressure at transmit frequency 4 MHz and 121 kPa peak-negative acoustic pressure ($n = 2$). Sensitivity of the subharmonic signal to 0–75 mmHg hydrostatic pressure was 0.15 dB/mmHg ($r^2 = 0.99$; $p < 0.001$).

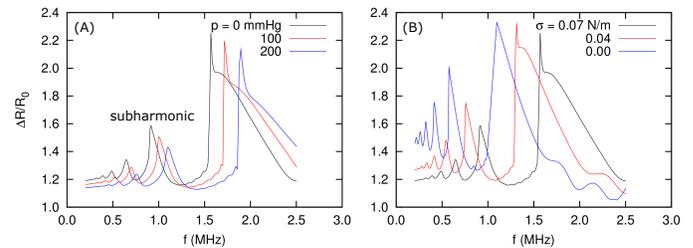


Fig. 6. Amplitude of bubble oscillation using the Rayleigh-Plesset equation with 150 kPa peak-to-peak acoustic pressure. (A) Simulation results at 0, 100 and 200 mmHg hydrostatic pressures. (B) Simulation results at and below the surface tension corresponding to air/water systems (0–0.07 N/m).

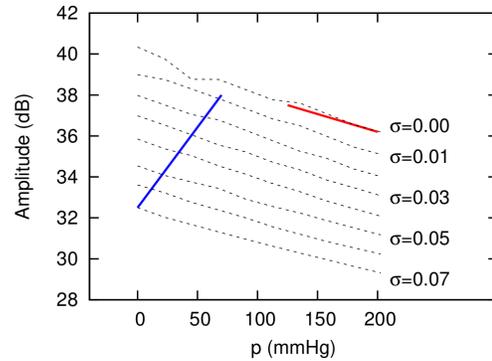


Fig. 7. Predicted subharmonic amplitude across 0–200 mmHg hydrostatic pressures and 0.00–0.07 N/m bubble surface tension, derived using the Rayleigh-Plesset equation with a bubble radius of 2 μm . The blue line shows a possible increase of the subharmonic signal with decreasing surface tension; the red line marks the decrease in subharmonic signal with increasing hydrostatic pressure at zero surface tension.

following exposure to 0–125 mmHg (Fig. 8). The greatest decrease in subharmonic amplitude between exposures to 125 and 200 mmHg hydrostatic pressures was 3.2 dB, which occurred at 181 kPa peak-negative acoustic pressure (26.6% maximum scanner acoustic output).

IV. DISCUSSION

In this study, we investigated the subharmonic response of the microbubble ultrasound contrast agent SonoVue at 0–200 mmHg hydrostatic pressures using a static pressure phantom.

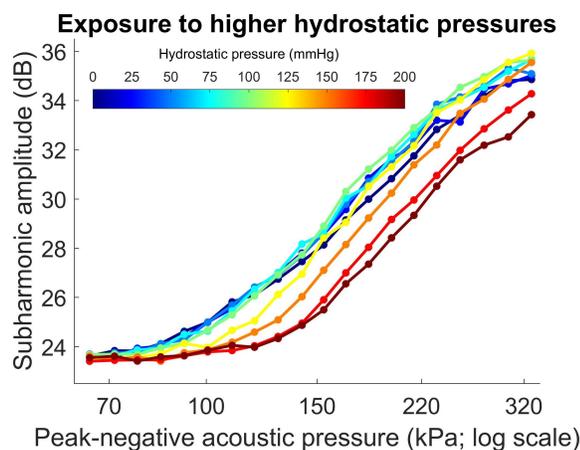


Fig. 8. Mean subharmonic amplitude of SonoVue at ambient hydrostatic pressure (0 mmHg), after 1 min of exposure to higher pressures ($n = 3$; standard deviations omitted for clarity). Data recorded with transmit frequency 5 MHz and 16-cycle pulses. Horizontal axis: water bath acoustic pressures corrected for 6.19 dB signal loss through the melamine foam and cassette window.

We developed a new phantom from commercially available components, an acquisition protocol to record radiofrequency data at incremental acoustic output levels, and a signal processing toolbox to extract the subharmonic amplitude from the received signal. The subharmonic signal of SonoVue increased from 0–50 mmHg hydrostatic pressure across 1.7–7.0 MHz transmit frequencies, and decreased from 125–200 mmHg. Decreasing surface tension may explain the increase in subharmonic signal from 0–50 mmHg pressure, while shell buckling and bubble destruction likely contribute to the decrease in subharmonic signal from 125–200 mmHg.

A. Optimal transmit frequency to use the subharmonic signal of SonoVue to estimate hydrostatic pressure

Across the transmit frequencies investigated, 4.0 MHz elicited the best subharmonic sensitivity for assessing hydrostatic pressures up to 75 mmHg (0.15 dB/mmHg). Comparable sensitivities to hydrostatic pressure were found for transmit frequencies 1.7–3.0 MHz, but the decrease in subharmonic signal from 75–125 mmHg would result in non-unique values that may be incorrectly interpreted when used to estimate pressures *in vivo*. This is not a concern at transmit frequency 4.0 MHz because the subharmonic signal plateaus from 75–125 mmHg. In addition, transmit frequencies 1.7–5.0 MHz involved extracting the subharmonic signal outside the bandwidth of the transducers used in this study. Future use of a transducer with a bandwidth that includes both the subharmonic and transmit frequencies may thus further increase upon the sensitivities found in this study. The ascending and descending pattern of the subharmonic-pressure relationship, nonetheless, are unique characteristics of SonoVue bubble behavior, and are independent of the choice of transducer.

B. The subharmonic-pressure response of SonoVue differs from other microbubble ultrasound contrast agents

The observed increase in the subharmonic signal of SonoVue from 0–75 mmHg hydrostatic pressure was in stark

contrast to our hypothesis predicting a linear decrease from 0–200 mmHg. Our hypothesis was based upon previous work that found this linear decrease in subharmonic signal across multiple ultrasound contrast agents — Sonazoid, Optison, Levovist and Definity — but SonoVue was not investigated in that study [10]. In agreement with our findings, Li and colleagues have recently reported an increase in the subharmonic amplitude of SonoVue from 0–50 mmHg hydrostatic pressure, followed by a decrease from 50–180 mmHg, at transmit frequency 4 MHz and peak-negative acoustic pressure 350 kPa [17]. At 450 kPa peak-negative acoustic pressure and even higher mechanical indices (≥ 0.225), however, Li and colleagues observed a decrease in the subharmonic amplitude of SonoVue with hydrostatic pressure [17]. This may be due to microbubble destruction at higher mechanical indices, and thus demonstrates the need for low mechanical index imaging in this new application of using SonoVue microbubbles to estimate hydrostatic pressure. A direct comparison of the subharmonic response of SonoVue with other microbubble ultrasound contrast agents will help verify this different behavior of SonoVue [31].

C. Decreasing surface tension may underpin the ascending phase of the SonoVue subharmonic-pressure relationship

Our simulations using the Rayleigh-Plesset equation revealed an increase in the subharmonic signal with decreasing bubble surface tension, and a decrease in the subharmonic signal with increasing hydrostatic pressure. These competing effects of hydrostatic pressure and surface tension likely underpin the resulting net subharmonic signal observed experimentally from SonoVue microbubbles. Based on our simulation results, we speculate that the observed increase in the subharmonic signal of SonoVue from 0–75 mmHg pressure was driven by a decreasing microbubble surface tension (Fig. 7). With decreasing surface tension, the oscillations of microbubbles shift toward more compression than expansion [11]. In addition to SonoVue microbubbles, which consist of sulfur hexafluoride gas (SF_6) in a phospholipid shell, this “compression-only” behavior has been implicated in the increased subharmonic signal of experimental perfluorobutane (C_4F_{10}) phospholipid-shell microbubbles with hydrostatic pressure [11].

D. Microbubble buckling and destruction contribute to the descending phase of the SonoVue subharmonic-pressure relationship

Above 125 mmHg hydrostatic pressure, we speculate that most of the SonoVue microbubbles are in a buckled state (i.e., zero surface tension), and thus the subharmonic signal of SonoVue decreases with increasing hydrostatic pressure (Fig. 7) [19]. The extent of the decrease of the subharmonic signal observed experimentally, however, was greater than that predicted by our simulations using the Rayleigh-Plesset equation. This discrepancy may be explained by irreversible bubble destruction at 150–200 mmHg pressures, which would violate the assumption that changes in void fraction solely reflect changes in bubble volume (v) — and not fewer bubbles

per unit volume (cf., Equation 2; $\frac{v_{\text{gas}}}{v_{\text{gas reference}}} = \frac{p_{\text{ref}}}{p_0}$). Indeed, the lower subharmonic signal observed in the hysteresis experiments following exposure to 150–200 mmHg pressures for 1 min, but not following exposure to 0–125 mmHg, support the inference of microbubble destruction above 125 mmHg. Destruction of SonoVue microbubbles upon exposure to higher hydrostatic pressures may include lipid shedding from the bubble shell and static diffusion of SF₆ gas out of the bubble core, followed by inertial cavitation and fragmentation of the bubble when the pressure is released [32]–[34].

Taken together, the descending phase of the subharmonic-pressure relationship of SonoVue from 125–200 mmHg is likely underpinned by both bubble buckling and bubble destruction. Our findings additionally reiterate the importance of empirical data in investigations of microbubble ultrasound contrast agents, as existing mathematical models do not yet fully characterize complex microbubble mechanics [12], [35].

E. Clinical implications

The subharmonic-pressure relationship of SonoVue may be used to estimate pressures *in vivo*. Whilst systolic and diastolic brachial artery pressures are measured routinely using a sphygmomanometer, blood pressures in the heart and aorta are currently assessed invasively using a pressure catheter. The linear increase in subharmonic amplitude from 0–75 mmHg may be applied to estimate left ventricular diastolic pressures (4–12 mmHg), which are critical for the assessment of diastolic performance [8], [9]. This pressure range is also appropriate for the assessment of right ventricular pressure (2–30 mmHg), and left (4–12 mmHg) and right atrial pressures (2–6 mmHg) across the cardiac cycle. On the other hand, the linear decrease in the subharmonic signal above 75 mmHg may be applied to assess aortic pressures across the cardiac cycle (80–120 mmHg). For clinical applications, an estimate of pressure within 5 mmHg of reference pressure is ideal [36], and appear to be most likely achievable with SonoVue at transmit frequencies 2.1, 4.0 and 5.0 MHz.

One possible obstacle in translating our ultrasound transmit and receive configurations to image the heart and great vessels could be a greater attenuation through the body than through the static pressure phantom in this study. However, the attenuation through 1 cm of septal or lateral myocardium at transmit frequency 2.1 MHz is less than 3 dB [37], [38], which is less than the measured 4.6–16.2 dB signal loss in our phantom. It is thus likely that the transmit and receive configurations developed in this study can be successfully translated to *in vivo* imaging. In addition, the transducers used in this study are already routinely used for cardiac and vascular imaging.

In addition to SonoVue, other microbubble ultrasound contrast agents such as Sonazoid and Definity have been investigated as potential pressure sensors. These have shown potential as non-invasive pressure sensors in the right ventricle [5], for diagnosing portal hypertension [39], and for estimating tumor interstitial fluid pressure in breast cancer [40]. One disadvantage of using SonoVue, compared with Sonazoid or Definity, is its non-unique subharmonic amplitudes from 0–200 mmHg hydrostatic pressures. However, an advantage of

SonoVue is its greater sensitivity to changes in pressure, as sensitivities of up to ± 0.16 dB/mmHg were found in this study. An earlier study by Halldorsdottir and colleagues comparing different microbubble ultrasound agents found lower sensitivities for Sonazoid (-0.08 dB/mmHg), Definity (-0.07 dB/mmHg) and Optison (-0.06 dB/mmHg), but did not investigate SonoVue [10]. Optimization of the ultrasound pulse shape subsequently improved the sensitivity of Sonazoid to -0.17 dB/mmHg for 0–40 mmHg hydrostatic pressures, but the maximum pressure range was limited by the decrease of the subharmonic signal below the noise floor [41]. Despite extensive research on Sonazoid, one of its limitations is that it is currently not approved for use in Europe [15]. The investigation of SonoVue, which is widely used in Europe, may help accelerate this exciting new technology from the lab to the clinic in Europe. Altogether, these characteristics need to be considered and balanced to ultimately choose the best possible microbubble ultrasound contrast agent for the target application in the human body.

F. Limitations and future work

This work included the development of a new static pressure phantom, suitable for microbubble experiments [20]. As the probe and cell culture cassette were positioned manually without the use of a rotation mount, the angle of incidence may not have been precisely 45°. Manual measurement with a protractor, however, allowed us to estimate the angle of incidence and with this, calculate a best-estimate of the signal loss prior to the pressure chamber. Due to resource limitations, our pressure meter was calibrated with a simple 2 m water column and only reached 147 mmHg hydrostatic pressure. To achieve a higher accuracy of pressures above 147 mmHg on the pressure meter, a pneumatic calibrator may be used in the future. In addition, only two sets of data were acquired at each hydrostatic pressure level and more sets would have been ideal. However, this was the result of a balance between minimizing the duration of experiments from the point of microbubble reconstitution (to minimize differences in microbubble properties over time), and maximizing the number of hydrostatic pressure levels ($n = 9$) and acoustic outputs investigated ($n = 20$). A visual inspection of the first two sets of data to determine any crossover of the datapoints was used instead to exclude erroneous datasets and to ensure the best estimate of the subharmonic amplitude of SonoVue in this study.

To provide first insights into the mechanisms underlying the subharmonic-pressure relationship of SonoVue, we used a model of a pure gas bubble with an effective surface tension [19], coupled with a simple representation of the bubble cloud response. This provided qualitative insight into our experimental findings, but does not fully represent the complex interaction between bubble clusters and acoustic waves [35]. Further work using more comprehensive and elaborate models, such as the subgrid model for bubbly cavitating flows proposed by Fuster and Colonius [42], will enable a better representation of the bubble cloud beyond commonly used single bubble models [12].

Building upon the ultrasound acquisition protocol and signal processing toolbox developed in this study, future work will

include shortening the protocol to be able to complete this test within 20 min in the clinic (the experimental protocol in this study investigating 0–200 mmHg hydrostatic pressures at one transmit frequency took up to 4.5 h to complete). This will likely include quicker identification of the optimal acoustic output and subsequent data collection across time at this single acoustic output [26], [43]. As differences in attenuation, blood viscosity [44] and temperature [45] likely affect the subharmonic signal of SonoVue, we envision that an incremental acoustic output scan will be necessary for each patient acquisition to identify an individualized optimal acoustic output, similar to previous work on Sonazoid [26] and Definity [5]. Studies conducted *in vitro* have found that an increase in temperature from room temperature to body temperature decreases microbubble stability [45], while an increase in viscosity from water to blood has the opposite effect of increasing microbubble stability [44]. Future experiments at body temperature and in a blood-mimicking fluid medium (instead of at room temperature and in water in this study) will therefore enable optimization of this technique prior to *in vivo* testing. In addition, amplitude modulation and pulse-shaping may be investigated as methods to further enhance the subharmonic signal of SonoVue [41], [46].

V. CONCLUSION

The subharmonic signal of SonoVue first increased with hydrostatic pressure across all experimental conditions (0–50 mmHg), and then decreased (125–200 mmHg). The increase in the subharmonic signal of SonoVue may be driven by a decreasing bubble surface tension, while the decrease may be attributed to both shell buckling and bubble destruction. Results report the largest sensitivity to date across 0–200 mmHg (± 0.16 dB/mmHg), opening promising translational perspectives for a less invasive method to assess diastolic filling pressures (compared with inserting a catheter into the heart).

APPENDIX

CHARACTERIZATION OF THE INSERTION LOSS THROUGH ONE WINDOW OF THE CLINICELL CHAMBER

Table V shows the insertion loss through one window of the CLINICell chamber at 0–50° angles of insonation, across 1–7 MHz transmit frequencies. Future ultrasound studies using the CLINICell cell culture cassette with 175 μm membrane can use this detailed characterization to calculate the acoustic signal loss in their experimental setups.

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REFERENCES

[1] M. Claudon, C. F. Dietrich, B. I. Choi, D. O. Cosgrove, M. Kudo, C. P. Nolsøe, F. Piscaglia, S. R. Wilson, R. G. Barr, M. C. Chammas *et al.*, “Guidelines and good clinical practice recommendations for contrast enhanced ultrasound (CEUS) in the liver – Update 2012,” *Ultrasound Med Biol*, vol. 39, no. 2, pp. 187–210, 2013.

[2] T. R. Porter, S. L. Mulvagh, S. S. Abdelmoneim, H. Becher, J. T. Belcik, M. Bierig, J. Choy, N. Gaibazzi, L. D. Gillam, R. Janardhanan *et al.*, “Clinical applications of ultrasonic enhancing agents in echocardiography: 2018 American Society of Echocardiography guidelines update,” *J Am Soc Echocardiogr*, vol. 31, no. 3, pp. 241–274, 2018.

[3] R. Senior, H. Becher, M. Monaghan, L. Agati, J. Zamorano, J. L. Vanoverschelde, P. Nihoyannopoulos, T. Edvardsen, and P. Lancellotti, “Clinical practice of contrast echocardiography: Recommendation by the European Association of Cardiovascular Imaging (EACVI) 2017,” *Eur Heart J Cardiovasc Imaging*, vol. 18, no. 11, pp. 1205–1205af, 2017.

[4] P. S. Sidhu, V. Cantisani, C. F. Dietrich, O. H. Gilja, A. Saftoiu, E. Bartels, M. Bertolotto, F. Calliada, D.-A. Clevert, D. Cosgrove *et al.*, “The EFSUMB guidelines and recommendations for the clinical practice of contrast-enhanced ultrasound (CEUS) in non-hepatic applications: Update 2017 (Long version),” *Ultraschall Med*, vol. 39, no. 02, pp. e2–e44, 2018.

[5] J. K. Dave, S. V. Kulkarni, P. P. Pangaonkar, M. Stanczak, M. E. McDonald, I. S. Cohen, P. Mehrotra, M. P. Savage, P. Walinsky, N. J. Ruggiero II *et al.*, “Non-invasive intra-cardiac pressure measurements using subharmonic-aided pressure estimation: Proof of concept in humans,” *Ultrasound Med Biol*, vol. 43, no. 11, pp. 2718–2724, 2017.

[6] I. Gupta, J. R. Eisenbrey, P. Machado, M. Stanczak, K. Wallace, and F. Forsberg, “On factors affecting subharmonic-aided pressure estimation (SHAPE),” *Ultrason Imaging*, vol. 41, no. 1, pp. 35–48, 2019.

[7] R. E. Klabunde, “Cardiovascular integration, adaptation, and pathophysiology,” in *Cardiovascular physiology concepts*, 2nd ed. Baltimore, MD: Lippincott Williams & Wilkins, 2012, ch. Cardiovascular integration, adaptation, and pathophysiology, pp. 198–234.

[8] J. Xi, P. Lamata, S. Niederer, S. Land, W. Shi, X. Zhuang, S. Ourselin, S. G. Duckett, A. K. Shetty, C. A. Rinaldi *et al.*, “The estimation of patient-specific cardiac diastolic functions from clinical measurements,” *Med Image Anal*, vol. 17, no. 2, pp. 133–146, 2013.

[9] J. Xi, W. Shi, D. Rueckert, R. Razavi, N. P. Smith, and P. Lamata, “Understanding the need of ventricular pressure for the estimation of diastolic biomarkers,” *Biomech Model Mechanobiol*, vol. 13, pp. 747–757, 2014.

[10] V. G. Halldorsdottir, J. K. Dave, L. M. Leodore, J. R. Eisenbrey, S. Park, A. L. Hall, K. Thomenius, and F. Forsberg, “Subharmonic contrast microbubble signals for noninvasive pressure estimation under static and dynamic flow conditions,” *Ultrason Imaging*, vol. 33, no. 3, pp. 153–164, 2011.

[11] P. J. Frinking, E. Gaud, J. Brochot, and M. Arditi, “Subharmonic scattering of phospholipid-shell microbubbles at low acoustic pressure amplitudes,” *IEEE Trans Ultrason Ferroelectr Freq Control*, vol. 57, no. 8, pp. 1762–1771, 2010.

[12] A. Katiyar, K. Sarkar, and F. Forsberg, “Modeling subharmonic response from contrast microbubbles as a function of ambient static pressure,” *J Acoust Soc Am*, vol. 129, no. 4, pp. 2325–2335, 2011.

[13] J. K. Dave, V. G. Halldorsdottir, J. R. Eisenbrey *et al.*, “On the implementation of an automated acoustic output optimization algorithm for subharmonic aided pressure estimation,” *Ultrasonics*, vol. 53, no. 4, pp. 880–888, 2013.

[14] J. Eisenbrey, J. Dave, V. Halldorsdottir, D. Merton, P. Machado, J. Liu, C. Miller, J. Gonzalez, S. Park, S. Dianis, C. Chalek, K. Thomenius, B. DB, V. Navarro, and F. Forsberg, “Simultaneous grayscale and subharmonic ultrasound imaging on a modified commercial scanner,” *Ultrasonics*, vol. 51, no. 8, pp. 890–897, 2011.

[15] A. W. Appis, M. J. Tracy, and S. B. Feinstein, “Update on the safety and efficacy of commercial ultrasound contrast agents in cardiac applications,” *Echo Res Pract*, vol. 2, no. 2, pp. R55–R62, 2015.

[16] K. S. Andersen and J. A. Jensen, “Impact of acoustic pressure on ambient pressure estimation using ultrasound contrast agent,” *Ultrasonics*, vol. 50, no. 2, pp. 294–299, 2010.

[17] F. Li, D. Li, and F. Yan, “Improvement of detection sensitivity of microbubbles as sensors to detect ambient pressure,” *Sensors*, vol. 18, no. 12, p. 4083, 2018.

[18] T. Sun, N. Jia, D. Zhang, and D. Xu, “Ambient pressure dependence of the ultra-harmonic response from contrast microbubbles,” *J Acoust Soc Am*, vol. 131, no. 6, pp. 4358–4364, 2012.

[19] P. Marmottant, S. van der Meer, M. Emmer, M. Versluis, N. de Jong, S. Hilgenfeldt, and D. Lohse, “A model for large amplitude oscillations of coated bubbles accounting for buckling and rupture,” *J Acoust Soc Am*, vol. 118, no. 6, pp. 3499–3505, 2005.

[20] A. Q. X. Nio, A. Faraci, K. Christensen-Jeffries, R. J. Eckersley, M. J. Monaghan, J. L. Raymond, F. Forsberg, and P. Lamata, “The subharmonic amplitude of SonoVue increases with hydrostatic pressure

TABLE V
CHARACTERIZATION OF THE INSERTION LOSS (DB) THROUGH ONE WINDOW OF THE CLINICELL CHAMBER. VALUES ARE MEAN (STANDARD DEVIATION) OF FOUR MEASUREMENTS.

Transmit frequency (MHz)	Angle (deg)										
	0	5	10	15	20	25	30	35	40	45	50
1.0	0.38 (0.08)	0.35 (0.09)	0.38 (0.08)	0.34 (0.10)	0.33 (0.01)	0.33 (0.11)	0.24 (0.03)	0.17 (0.06)	0.43 (0.05)	0.92 (0.05)	1.91 (0.20)
1.5	0.68 (0.02)	0.69 (0.02)	0.69 (0.04)	0.75 (0.05)	0.78 (0.03)	0.86 (0.03)	0.96 (0.02)	1.01 (0.01)	1.33 (0.04)	2.04 (0.06)	4.28 (0.12)
2.0	1.11 (0.02)	1.13 (0.01)	1.20 (0.00)	1.29 (0.02)	1.43 (0.03)	1.61 (0.03)	1.84 (0.03)	2.12 (0.02)	2.79 (0.04)	4.36 (0.08)	8.84 (0.17)
2.5	1.41 (0.03)	1.72 (0.02)	2.38 (0.03)	3.06 (0.04)	3.58 (0.01)	3.94 (0.04)	4.39 (0.05)	5.02 (0.06)	6.50 (0.04)	10.45 (0.28)	24.25 (0.74)
3.0	1.53 (0.02)	1.46 (0.02)	1.33 (0.04)	1.37 (0.05)	1.95 (0.04)	4.02 (0.07)	7.96 (0.10)	12.16 (0.17)	15.64 (0.17)	14.30 (0.39)	7.92 (0.29)
3.5	1.52 (0.02)	1.48 (0.03)	1.37 (0.02)	1.24 (0.03)	1.12 (0.02)	1.18 (0.05)	1.70 (0.07)	2.82 (0.04)	4.07 (0.03)	4.34 (0.11)	3.51 (0.16)
4.0	1.42 (0.02)	1.40 (0.03)	1.31 (0.02)	1.23 (0.03)	1.14 (0.00)	1.11 (0.01)	1.26 (0.02)	1.60 (0.04)	2.13 (0.08)	2.63 (0.05)	2.99 (0.02)
4.5	1.20 (0.02)	1.19 (0.02)	1.13 (0.01)	1.08 (0.02)	1.04 (0.00)	1.09 (0.00)	1.23 (0.01)	1.51 (0.05)	1.99 (0.08)	2.72 (0.09)	3.44 (0.05)
5.0	0.96 (0.01)	1.01 (0.02)	1.19 (0.02)	1.44 (0.02)	1.64 (0.01)	1.84 (0.01)	1.95 (0.01)	2.02 (0.08)	2.35 (0.16)	3.22 (0.17)	3.93 (0.07)
5.5	0.83 (0.01)	1.32 (0.06)	2.67 (0.09)	4.50 (0.11)	6.20 (0.10)	7.43 (0.07)	7.70 (0.02)	6.19 (0.14)	3.40 (0.10)	3.64 (0.14)	4.05 (0.02)
6.0	0.65 (0.01)	0.78 (0.05)	1.18 (0.07)	1.83 (0.10)	2.64 (0.13)	3.45 (0.11)	4.00 (0.09)	4.00 (0.07)	3.53 (0.13)	4.25 (0.03)	3.96 (0.07)
6.5	0.68 (0.01)	0.72 (0.04)	0.87 (0.06)	1.15 (0.09)	1.62 (0.11)	2.25 (0.11)	2.80 (0.10)	3.13 (0.13)	3.71 (0.23)	7.23 (0.16)	4.89 (0.34)
7.0	0.88 (0.01)	0.86 (0.05)	0.87 (0.05)	0.95 (0.08)	1.23 (0.13)	1.73 (0.15)	2.30 (0.12)	2.93 (0.20)	4.27 (0.38)	10.50 (0.57)	7.21 (0.75)

at low incident acoustic pressures,” in *Ultrasonics Symposium (IUS), 2017 IEEE International*, 2017.

[21] I. Beekers, T. van Rooij, A. F. van der Steen, N. de Jong, M. D. Verweij, and K. Kooiman, “Acoustic characterization of the CLINICell for ultrasound contrast agent studies,” *IEEE Trans Ultrason Ferroelectr Freq Control*, vol. 66, no. 1, pp. 244–246, 2019.

[22] V. Halldorsdottir, J. Dave, J. Eisenbrey, P. Machado, H. Zhao, J. Liu, D. Merton, and F. Forsberg, “Subharmonic aided pressure estimation for monitoring interstitial fluid pressure in tumours – in vitro and in vivo proof of concept,” *Ultrasonics*, vol. 54, no. 7, pp. 1938–1944, 2014.

[23] J. L. Raymond, K. J. Haworth, K. B. Bader, K. Radhakrishnan, J. K. Griffin, S.-L. Huang, D. D. McPherson, and C. K. Holland, “Broadband attenuation measurements of phospholipid-shelled ultrasound contrast agents,” *Ultrasound Med Biol*, vol. 40, no. 2, pp. 410–421, 2014.

[24] W. T. Shi, F. Forsberg, J. S. Raichlen, L. Needleman, and B. B. Goldberg, “Pressure dependence of subharmonic signals from contrast microbubbles,” *Ultrasound Med Biol*, vol. 25, no. 2, pp. 275–283, 1999.

[25] J. K. Dave, V. G. Halldorsdottir, J. R. Eisenbrey, and F. Forsberg, “Processing of subharmonic signals from ultrasound contrast agents to determine ambient pressures,” *Ultrason Imaging*, vol. 34, no. 2, pp. 81–92, 2012.

[26] J. K. Dave, V. G. Halldorsdottir, J. R. Eisenbrey, J. S. Raichlen, J.-B. Liu, M. E. McDonald, K. Dickie, S. Wang, C. Leung, and F. Forsberg, “Noninvasive LV pressure estimation using subharmonic emissions from microbubbles,” *JACC Cardiovasc Imaging*, vol. 5, no. 1, pp. 87–92, 2012.

[27] L. Bergamasco and D. Fuster, “Oscillation regimes of gas/vapor bubbles,” *Int J Heat Mass Transfer*, vol. 112, pp. 72–80, 2017.

[28] C. Greis, “Technology overview: SonoVue (Bracco, Milan),” *Eur Radiol*, vol. 14, no. 8 Supplement, pp. 11–15, 2004.

[29] A. Prosperetti, “Subharmonics and ultraharmonics in the forced oscillations of weakly nonlinear systems,” *Am J Phys*, vol. 44, no. 6, pp. 548–554, 1976.

[30] E. Carstensen and L. Foldy, “Propagation of sound through a liquid containing bubbles,” *J Acoust Soc Am*, vol. 19, no. 3, pp. 481–501, 1947.

[31] I. Gupta, A. Q. X. Nio, A. Faraci, M. Torkzaban, K. Christensen-Jeffries, K. Nam, J. L. Raymond, K. Wallace, M. J. Monaghan, D. Fuster, R. J. Eckersley, P. Lamata, and F. Forsberg, “The effects of hydrostatic pressure on the subharmonic response of SonoVue and Sonazoid,” in *Ultrasonics Symposium (IUS), 2019 IEEE International*, 2019, In Press.

[32] J. E. Chomas, P. Dayton, J. Allen, K. Morgan, and K. W. Ferrara, “Mechanisms of contrast agent destruction,” *IEEE Trans Ultrason Ferroelectr Freq Control*, vol. 48, no. 1, pp. 232–248, 2001.

[33] E. Kanbar, D. Fouan, C. A. Sennoga, A. A. Doinikov, and A. Bouakaz, “Impact of filling gas on subharmonic emissions of phospholipid ultrasound contrast agents,” *Ultrasound Med Biol*, vol. 43, no. 5, pp. 1004–1015, 2017.

[34] Y. Luan, G. Lajoinie, E. Gelderblom, I. Skachkov, A. F. van der Steen, H. J. Vos, M. Versluis, and N. De Jong, “Lipid shedding from single oscillating microbubbles,” *Ultrasound Med Biol*, vol. 40, no. 8, pp. 1834–1846, 2014.

[35] D. Fuster, “A review of models for bubble clusters in cavitating flows,” *Flow Turbul Combust*, pp. 1–40, 2018.

[36] G. S. Stergiou, B. Alpert, S. Mieke, R. Asmar, N. Atkins, S. Eckert, G. Frick, B. Friedman, T. Graßl, T. Ichikawa *et al.*, “A universal standard for the validation of blood pressure measuring devices: Association for the Advancement of Medical Instrumentation/European Society of Hypertension/International Organization for Standardization (AAMI/ESH/ISO) collaboration statement,” *Hypertension*, vol. 71, no. 3, pp. 368–374, 2018.

[37] A. Q. X. Nio, E. J. Stöhr, and R. Shave, “Age-related differences in left ventricular structure and function between healthy men and women,” *Climacteric*, vol. 20, no. 5, pp. 476–483, 2017.

[38] D. E. Sosnovik, S. L. Baldwin, S. H. Lewis, M. R. Holland, and J. G. Miller, “Transmural variation of myocardial attenuation measured with a clinical imager,” *Ultrasound Med Biol*, vol. 27, no. 12, pp. 1643–1650, 2001.

[39] J. R. Eisenbrey, J. K. Dave, V. G. Halldorsdottir, D. A. Merton, C. Miller, J. M. Gonzalez, P. Machado, S. Park, S. Dianis, C. L. Chalek, C. E. Kim, J. P. Baliff, K. E. Thomenis, D. B. Brown, V. Navarro, and F. Forsberg, “Chronic liver disease: Noninvasive subharmonic aided pressure estimation of hepatic venous pressure gradient,” *Radiology*, vol. 268, no. 2, pp. 581–588, 2013.

[40] K. Nam, J. R. Eisenbrey, M. Stanczak, A. Sridharan, A. C. Berger, T. Avery, J. P. Palazzo, and F. Forsberg, “Monitoring neoadjuvant chemotherapy for breast cancer by using three-dimensional subharmonic aided pressure estimation and imaging with us contrast agents: Preliminary experience,” *Radiology*, vol. 285, no. 1, pp. 53–62, 2017.

[41] I. Gupta, J. Eisenbrey, M. Stanczak, A. Sridharan, J. K. Dave, J.-B. Liu, C. Hazard, X. Wang, P. Wang, H. Li, K. Wallace, and F. Forsberg, “Effect of pulse shaping on subharmonic aided pressure estimation in vitro and in vivo,” *J Ultrasound Med*, vol. 36, no. 1, pp. 3–11, 2017.

[42] D. Fuster and T. Colonius, “Modelling bubble clusters in compressible liquids,” *J Fluid Mech*, vol. 688, pp. 352–389, 2011.

[43] C. Esposito, K. Dickie, F. Forsberg, and J. K. Dave, “Towards real-time implementation of subharmonic aided pressure estimation (SHAPE)–How to identify optimum acoustic output for SHAPE?” in *Ultrasonics Symposium (IUS), 2017 IEEE International*, 2017.

[44] B. Helfield, J. J. Black, B. Qin, J. Pacella, X. Chen, and F. S. Villanueva, “Fluid viscosity affects the fragmentation and inertial cavitation threshold of lipid-encapsulated microbubbles,” *Ultrasound Med Biol*, vol. 42, no. 3, pp. 782–794, 2016.

[45] H. Mulvana, E. Stride, J. V. Hajnal, and R. J. Eckersley, “Temperature dependent behavior of ultrasound contrast agents,” *Ultrasound Med Biol*, vol. 36, no. 6, pp. 925–934, 2010.

[46] R. J. Eckersley, C. T. Chin, and P. N. Burns, “Optimising phase and amplitude modulation schemes for imaging microbubble contrast agents at low acoustic power,” *Ultrasound Med Biol*, vol. 31, no. 2, pp. 213–219, 2005.



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